

APPENDIX A

Some Useful Mathematical Formulas

A.1 Useful Vector Identities

$$\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = \mathbf{b} \cdot (\mathbf{c} \times \mathbf{a}) = \mathbf{c} \cdot (\mathbf{a} \times \mathbf{b}), \quad (\text{A.1})$$

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = \mathbf{b}(\mathbf{a} \cdot \mathbf{c}) - \mathbf{c}(\mathbf{a} \cdot \mathbf{b}), \quad (\text{A.2})$$

$$\nabla \times \nabla \psi = 0, \quad (\text{A.3})$$

$$\nabla \cdot \nabla \times \mathbf{A} = 0, \quad (\text{a.4})$$

$$\nabla \cdot (\psi \mathbf{A}) = \mathbf{A} \cdot \nabla \psi + \psi \nabla \cdot \mathbf{A}, \quad (\text{A.5})$$

$$\nabla \times (\psi \mathbf{A}) = \nabla \psi \times \mathbf{A} + \psi \nabla \times \mathbf{A}, \quad (\text{A.6})$$

$$\nabla \cdot (\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot \nabla \times \mathbf{A} - \mathbf{A} \cdot \nabla \times \mathbf{B}, \quad (\text{A.7})$$

$$\nabla(\mathbf{A} \cdot \mathbf{B}) = (\mathbf{A} \cdot \nabla)\mathbf{B} + (\mathbf{B} \cdot \nabla)\mathbf{A} + \mathbf{A} \times \nabla \times \mathbf{B} + \mathbf{B} \times \nabla \times \mathbf{A}, \quad (\text{A.8})$$

$$\nabla \times (\mathbf{A} \times \mathbf{B}) = (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} + \mathbf{A} \nabla \cdot \mathbf{B} - \mathbf{B} \nabla \cdot \mathbf{A}, \quad (\text{A.9})$$

$$\nabla \times \nabla \times \mathbf{A} = \nabla \nabla \cdot \mathbf{A} - \nabla^2 \mathbf{A}. \quad (\text{A.10})$$

In Cartesian coordinates, $\nabla^2 \mathbf{A}$ can be decomposed as

$$\nabla^2 \mathbf{A} = \hat{x} \nabla^2 A_x + \hat{y} \nabla^2 A_y + \hat{z} \nabla^2 A_z, \quad (\text{A.11})$$

because ∇^2 commutes with \hat{x} , \hat{y} , and \hat{z} , i.e., $\nabla^2 \hat{x} = \hat{x} \nabla^2$ and so on. This is not true in other curvilinear coordinates; hence, this decomposition is not allowed.

A.2 Gradient, Divergence, Curl, and Laplacian in Rectangular, Cylindrical, Spherical, and General Orthogonal Curvilinear Coordinate Systems

(a) Rectangular System; x, y, z :

$$\nabla \psi = \frac{\partial \psi}{\partial x} \hat{x} + \frac{\partial \psi}{\partial y} \hat{y} + \frac{\partial \psi}{\partial z} \hat{z}, \quad (\text{A.12})$$

$$\nabla \cdot \mathbf{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}, \quad (\text{A.13})$$

$$\nabla \times \mathbf{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{x} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{y} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{z}, \quad (\text{A.14})$$

$$\nabla^2 \psi = \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2}. \quad (\text{A.15})$$

(b) **Cylindrical System; ρ, ϕ, z :**

$$\nabla\psi = \frac{\partial\psi}{\partial\rho}\hat{\rho} + \frac{1}{\rho}\frac{\partial\psi}{\partial\phi}\hat{\phi} + \frac{\partial\psi}{\partial z}\hat{z}, \quad (\text{A.16})$$

$$\nabla \cdot \mathbf{A} = \frac{1}{\rho}\frac{\partial}{\partial\rho}(\rho A_\rho) + \frac{1}{\rho}\frac{\partial A_\phi}{\partial\phi} + \frac{\partial A_z}{\partial z}, \quad (\text{A.17})$$

$$\nabla \times \mathbf{A} = \left(\frac{1}{\rho}\frac{\partial A_z}{\partial\phi} - \frac{\partial A_\phi}{\partial z}\right)\hat{\rho} + \left(\frac{\partial A_\rho}{\partial z} - \frac{\partial A_z}{\partial\rho}\right)\hat{\phi} + \frac{1}{\rho}\left(\frac{\partial}{\partial\rho}(\rho A_\phi) - \frac{\partial A_\rho}{\partial\phi}\right)\hat{z}, \quad (\text{A.18})$$

$$\nabla^2\psi = \frac{1}{\rho}\frac{\partial}{\partial\rho}\left(\rho\frac{\partial\psi}{\partial\rho}\right) + \frac{1}{\rho^2}\frac{\partial^2\psi}{\partial\phi^2} + \frac{\partial^2\psi}{\partial z^2}. \quad (\text{A.19})$$

(c) **Spherical System; r, θ, ϕ :**

$$\nabla\psi = \frac{\partial\psi}{\partial r}\hat{r} + \frac{1}{r}\frac{\partial\psi}{\partial\theta}\hat{\theta} + \frac{1}{r\sin\theta}\frac{\partial\psi}{\partial\phi}\hat{\phi}, \quad (\text{A.20})$$

$$\nabla \cdot \mathbf{A} = \frac{1}{r^2}\frac{\partial}{\partial r}(r^2 A_r) + \frac{1}{r\sin\theta}\frac{\partial}{\partial\theta}(\sin\theta A_\theta) + \frac{1}{r\sin\theta}\frac{\partial A_\phi}{\partial\phi}, \quad (\text{A.21})$$

$$\begin{aligned} \nabla \times \mathbf{A} = \frac{1}{r\sin\theta} \left[\frac{\partial}{\partial\theta}(\sin\theta A_\phi) - \frac{\partial A_\theta}{\partial\phi} \right] \hat{r} + \frac{1}{r} \left[\frac{1}{\sin\theta}\frac{\partial A_r}{\partial\phi} - \frac{\partial}{\partial r}(r A_\phi) \right] \hat{\theta} \\ + \frac{1}{r} \left[\frac{\partial}{\partial r}(r A_\theta) - \frac{\partial A_r}{\partial\theta} \right] \hat{\phi}, \end{aligned} \quad (\text{A.22})$$

$$\nabla^2\psi = \frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial\psi}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial\psi}{\partial\theta}\right) + \frac{1}{r^2\sin^2\theta}\frac{\partial^2\psi}{\partial\phi^2}. \quad (\text{A.23})$$

(d) **General Orthogonal Curvilinear Coordinate System; x_1, x_2, x_3 :**

The metric coefficients (h_1, h_2, h_3) in a general orthogonal curvilinear coordinate system are defined by

$$ds_i = h_i dx_i; \quad i = 1 \text{ or } 2, \text{ or } 3, \quad (\text{A.24})$$

where ds_i denotes a differential length in the direction of dx_i . Moreover, the variable, x_i may not have the dimension of length. One way of finding the metric coefficients is to express the rectangular variables in terms of the variables of that system:

$$x = x(x_1, x_2, x_3),$$

$$y = y(x_1, x_2, x_3),$$

$$z = z(x_1, x_2, x_3).$$

Then

$$ds_i = \left[\left(\frac{\partial x}{\partial x_i} \right)^2 + \left(\frac{\partial y}{\partial x_i} \right)^2 + \left(\frac{\partial z}{\partial x_i} \right)^2 \right]^{1/2} dx_i, \quad i = 1, 2, 3. \quad (\text{A.25})$$

Hence,

$$h_i = \left[\left(\frac{\partial x}{\partial x_i} \right)^2 + \left(\frac{\partial y}{\partial x_i} \right)^2 + \left(\frac{\partial z}{\partial x_i} \right)^2 \right]^{1/2}. \quad (\text{A.26})$$

For instance, in an elliptical coordinate system,

$$x = c \cosh u \cos v, \quad (\text{A.27})$$

$$y = c \sinh u \sin v. \quad (\text{A.28})$$

If (x_1, x_2, x_3) represent (u, v, z) , then by applying (26), we have

$$h_1 = h_2 = c(\sinh^2 u \cos^2 v + \cosh^2 u \sin^2 v)^{1/2} = c(\cosh^2 u - \cos^2 v)^{1/2}, \quad (\text{A.29})$$

$$h_3 = 1. \quad (\text{A.30})$$

In general, for any orthogonal curvilinear coordinate system,

$$\nabla \psi = \sum_{i=1}^3 \frac{1}{h_i} \frac{\partial \psi}{\partial x_i} \hat{x}_i, \quad (\text{A.31})$$

$$\nabla \cdot \mathbf{A} = \frac{1}{\Delta} \sum_{i=1}^3 \frac{\partial}{\partial x_i} \left(\frac{\Delta A_i}{h_i} \right), \quad \Delta = h_1 h_2 h_3, \quad (\text{A.32})$$

$$\nabla \times \mathbf{A} = \frac{1}{\Delta} \begin{vmatrix} h_1 \hat{x}_1 & h_2 \hat{x}_2 & h_3 \hat{x}_3 \\ \frac{\partial}{\partial x_1} & \frac{\partial}{\partial x_2} & \frac{\partial}{\partial x_3} \\ h_1 A_1 & h_2 A_2 & h_3 A_3 \end{vmatrix}, \quad (\text{A.33})$$

$$\nabla^2 \psi = \frac{1}{\Delta} \sum_{i=1}^3 \frac{\partial}{\partial x_i} \left(\frac{\Delta}{h_i^2} \frac{\partial \psi}{\partial x_i} \right). \quad (\text{A.34})$$

A.3 Useful Integral Identities

In the following formulas, V is a volume bounded by a closed surface S . The unit vector \hat{n} is normal to S and points outward.

(a) **Gradient Identity:**

$$\oint_V \nabla \phi dV = \oint_S \phi \hat{n} dS. \quad (\text{A.35})$$

(b) Gauss' Divergence Theorem:

$$\oint_V \nabla \cdot \mathbf{A} dV = \oint_S \mathbf{A} \cdot \hat{n} dS. \quad (\text{A.36})$$

(c) Vector Stokes' Theorem:

$$\oint_V \nabla \times \mathbf{A} dV = \oint_S \hat{n} \times \mathbf{A} dS. \quad (\text{A.37})$$

(d) First Form of Green's Theorem:

$$\oint_V [\phi_1 \nabla^2 \phi_2 + \nabla \phi_1 \cdot \nabla \phi_2] dV = \oint_S \hat{n} \cdot \phi_1 \nabla \phi_2 dS. \quad (\text{A.38})$$

(e) Second Form of Green's Theorem:

$$\oint_V [\phi_1 \nabla^2 \phi_2 - \phi_2 \nabla^2 \phi_1] dV = \oint_S \hat{n} \cdot (\phi_1 \nabla \phi_2 - \phi_2 \nabla \phi_1) dS. \quad (\text{A.39})$$

(f) Vector Green's Theorem:

$$\begin{aligned} \oint_V [\mathbf{P} \cdot \nabla \times \nabla \times \mathbf{Q} - \mathbf{Q} \cdot \nabla \times \nabla \times \mathbf{P}] dV \\ = \oint_S [\mathbf{Q} \times \nabla \times \mathbf{P} - \mathbf{P} \times \nabla \times \mathbf{Q}] \cdot \hat{n} dS. \end{aligned} \quad (\text{A.40})$$

The above may all be proved from Gauss' divergence theorem.

(g) Stokes' Theorem:

If S is an unclosed surface bounded by a contour C , then

$$\int_S (\nabla \times \mathbf{A}) \cdot \hat{n} dS = \oint_C \mathbf{A} \cdot d\mathbf{l}, \quad (\text{A.41})$$

$$\int_S \hat{n} \times \nabla \phi dS = \oint_C \phi d\mathbf{l}. \quad (\text{A.42})$$

(h) Gauss' Theorem in Two Dimensions:

$$\int_S (\nabla \cdot \mathbf{A}) dS = \oint_C \mathbf{A} \cdot \hat{n} dl, \quad (\text{A.43})$$

The above identities for tensors and dyads can also be readily established (see Appendix B).

A.4 Integral Transforms**(a) Fourier:**

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dy e^{ixy} \tilde{f}(y), \quad (\text{A.44})$$

$$\tilde{f}(y) = \int_{-\infty}^{\infty} dx e^{-ixy} f(x), \quad (\text{A.45})$$

$$\delta(x - x') = \frac{1}{2\pi} \int_{-\infty}^{\infty} dy e^{i(x-x')y}. \quad (\text{A.46})$$

(b) Cylindrical Hankel:

$$f(\rho) = \int_0^{\infty} d\lambda \lambda J_n(\lambda\rho) \tilde{f}(\lambda), \quad (\text{A.47})$$

$$\tilde{f}(\lambda) = \int_0^{\infty} d\rho \rho J_n(\lambda\rho) f(\rho), \quad (\text{A.48})$$

$$\frac{\delta(\rho - \rho')}{\rho} = \int_0^{\infty} d\lambda \lambda J_n(\lambda\rho) J_n(\lambda\rho'), \quad (\text{A.49})$$

where $J_n(x)$ is a cylindrical Bessel function of n -th order.

(c) Spherical Hankel:

$$f(r) = \left(\frac{2}{\pi}\right)^{1/2} \int_0^{\infty} d\lambda \lambda^2 j_n(\lambda r) \tilde{f}(\lambda), \quad (\text{A.50})$$

$$\tilde{f}(\lambda) = \left(\frac{2}{\pi}\right)^{1/2} \int_0^{\infty} dr r^2 j_n(\lambda r) f(r), \quad (\text{A.51})$$

$$\frac{\delta(r - r')}{r^2} = \frac{2}{\pi} \int_0^{\infty} d\lambda \lambda^2 j_n(\lambda r) j_n(\lambda r'), \quad (\text{A.52})$$

where $j_n(x)$ is a spherical Bessel function of n -th order.

(d) Hilbert:

$$g(t) = \frac{1}{\pi} P.V. \int_{-\infty}^{\infty} d\tau \frac{f(\tau)}{\tau - t}, \quad (\text{A.53})$$

$$f(\tau) = -\frac{1}{\pi} P.V. \int_{-\infty}^{\infty} dt \frac{g(t)}{t - \tau}, \quad (\text{A.54})$$

$$\delta(t - t') = -\frac{1}{\pi^2} P.V. \int_{-\infty}^{\infty} d\tau \frac{1}{\tau - t} P.V. \frac{1}{t' - \tau}, \quad (\text{A.55})$$

where $P.V.\frac{1}{x}$ is regarded as a generalized function (see Appendix C).

(e) Radon:

A Radon transform in an n -dimensional space is defined as

$$\tilde{f}(p, \hat{\xi}) = \int_{-\infty}^{\infty} f(\mathbf{x}) \delta(p - \hat{\xi} \cdot \mathbf{x}) d\mathbf{x}, \quad (\text{A.56})$$

where \mathbf{x} is a position vector in an n -dimensional space, $\hat{\xi}$ is a unit vector, and $d\mathbf{x}$ implies a volume integral in an n -dimensional space. The inverse Radon transform is different in even and odd dimensions. In even dimensions, it is

$$f(\mathbf{x}) = \frac{1}{(2\pi i)^n} \int_{|\xi|=1} d\xi \int_{-\infty}^{\infty} dp \frac{1}{p - \hat{\xi} \cdot \mathbf{x}} \left(\frac{\partial}{\partial p} \right)^{n-1} \tilde{f}(p, \hat{\xi}), \quad (\text{A.57})$$

where the $d\xi$ integral implies integrating over all angles of $\hat{\xi}$, i.e., $d\xi$ is an elemental area on the surface of a unit sphere. In odd dimensions,

$$f(\mathbf{x}) = \frac{1}{2} \frac{1}{(2\pi i)^{n-1}} \int_{|\xi|=1} d\xi \left(\frac{\partial}{\partial p} \right)^{n-1} \tilde{f}(p, \hat{\xi}) \Big|_{p=\hat{\xi} \cdot \mathbf{x}}. \quad (\text{A.58})$$